LARGE-TIME BEHAVIOR OF SOLUTIONS TO A SCALAR CONSERVATION LAW IN SEVERAL SPACE DIMENSIONS

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ABSTRACT. We consider solutions of the Cauchy problem in \mathbf{R}_+^{n+1} for the equation $u_t + \operatorname{div}_x f(u) = 0$. The initial data is assumed to be a compact perturbation of a function of the form, $\varphi(x) = a$ for $\langle x, \mu \rangle > 0$, $\varphi(x) = b$ for $\langle x, \mu \rangle < 0$, where a and b are constants and μ is a given unit vector. The Cauchy problem together with an entropy condition on u is known to be well posed. The solution with unperturbed initial data, $\varphi(x)$, is a traveling shock, $\varphi(x-\vec{k}t)$, provided that $\varphi(x-\vec{k}t)$ satisfies the entropy condition (an inequality on a, b, μ , and f). Assuming this type of condition on φ , we study the large-time behavior of u. In particular, we show that u converges to a traveling shock whose profile agrees with $\langle x, \mu \rangle = 0$ outside of a compact set.

Introduction. In 1958 A. M. Ilin and O. A. Oleinik [4] studied the large-time stability of a shock solution to a single conservation law. They considered the Cauchy problem

(0.1)
$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x} f(u) = 0 \text{ for } t > 0, \ x \in \mathbf{R},$$

(0.2)
$$u(0,x) = u_0(x)$$
 for $x \in \mathbf{R}$,

where u_0 is bounded,

$$u_0(x) = a \cdot \chi_{\mathbf{R}^+}(x) + b \cdot \chi_{\mathbf{R}^-}(x) \equiv \psi(x)$$

for |x| sufficiently large,

$$(0.3) a < b and f'' > 0.$$

Solutions of the Cauchy problem are unique among those which satisfy an entropy condition (see (E) in §1). The structure assumptions, (0.3), ensure that $\psi(x-kt)$ is a solution of (0.1) and satisfies the entropy condition when k=(f(b)-f(a))/(b-a). Ilin and Oleinik used the viscosity method to show that for all t sufficiently large, $u(t,x)=\psi(x_0+x-kt)$ for some $x_0\in\mathbf{R}$. Thus if the initial data is a compact perturbation of ψ , the solution evolves after finite time into the traveling shock solution, $\psi(x_0+x-kt)$. We remark that the same result was proved by C. Dafermos [2] and T. P. Liu [7] using the notion of generalized characteristic curves.

In this paper we investigate the n-dimensional analog of this problem:

(0.4)
$$\partial u/\partial t + \operatorname{div} f(u) = 0 \text{ for } t > 0, \ x \in \mathbf{R}^n,$$

$$(0.5) u(0,x) = u_0(x) for x \in \mathbf{R}^n$$

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where $f \in C^1(\mathbf{R}; \mathbf{R}^n)$ and $n \ge 1$. We assume that $u_0 \in L^{\infty}(\mathbf{R}^n)$ and $u_0 = \varphi$ for $|x| \ge R_1 > 0$, where

$$\varphi(x) = \left\{ \begin{array}{ll} a & \text{for } \langle x, \mu \rangle > 0, \\ b & \text{for } \langle x, \mu \rangle < 0 \end{array} \right.$$

for some unit vector $\mu \in \mathbf{R}^n$ and a < b. We define a solution of (0.4) to be a distribution solution in $C([0,\infty); L^1_{loc}(\mathbf{R}^n))$ which satisfies Vol'pert's entropy condition, (E). (See §1.)

If k = [f(b) - f(a)]/(b-a), it follows from (E) that $\varphi(x - kt)$ is a solution of (0.4) if and only if $\langle \tilde{f}(c), \mu \rangle \leq 0$ for $a \leq c \leq b$ where $\tilde{f}(c) \equiv f(c) - f(a) - (c-a)k$. We assume throughout the paper that φ is nondegenerate, i.e., the above entropy condition holds in the following strict sense: there exists a constant $\theta > 0$ such that

(0.6)
$$\langle \mu, \tilde{f}(c) \rangle \leq -\theta(b-c)(c-a) \text{ for } a \leq c \leq b.$$

We summarize our results as follows. In §§1 and 2 we show that there are bounds $m_1 < a$ and $m_2 > b$ (determined explicitly by \tilde{f}) so that if

$$(0.7) m_1 < \operatorname{ess inf} u_0 \quad \text{and} \quad m_2 > \operatorname{ess sup} u_0,$$

then after a finite time, $a \leq u(t,x) \leq b$. Moreover, there is a set $M \subset \mathbb{R}^n$ and a constant $R_2 > 0$ such that if

$$v(x) = b \cdot \chi_{\mathbf{M}}(X) + a \cdot \chi_{\mathbf{R}^n - \mathbf{M}}(x),$$

then

$$u(t,x) - v(x-kt) \to 0$$
 in $L^1(\mathbf{R}^n)$ as $t \to \infty$

and

$$M \cap \{x: |x| > R_2\} = \{x: \langle x, \mu \rangle < 0, |x| > R_2\}.$$

Our investigation shows that the large-time behavior of u can be described more definitely and that it depends on $d \equiv$ dimension (span $\{\tilde{f}(c): a \leq c \leq b\}$). In §3 we study the case d = n and prove the following results. Here the initial data is called admissible if it satisfies (0.7) and $u_0 = \varphi$ for $|x| \geq R_1 > 0$.

THEOREM 3.2. If u_0 is admissible and d=n, then M is a Lipschitz domain. More precisely, there is a set of coordinates $(y_1, \ldots, y_n) = (\hat{y}, y_n)$ obtained by a rotation and a constant, C, (both depending only on \tilde{f}) so that $M = \{(\hat{y}, y_n): y_n < g(\hat{y})\}$ where g is Lipschitz continuous and $||\nabla g||_{\infty} \leq C$.

We prove the following stability estimate for the shock front, g.

THEOREM 3.3. Suppose u_1 and u_2 are admissible initial data and d = n. Let g_1 and g_2 be the end state shock fronts. Then

$$\sup_{\hat{y} \in \mathbf{R}^{n-1}} |g_1(\hat{y}) - g_2(\hat{y})| \le C(\tilde{f}) \cdot \left(\int_{\mathbf{R}^n} |u_1 - u_2| dx \right)^{1/n}.$$

Using the above estimate, we show that in finite time, u is equal to the traveling shock, v(x-kt), outside an arbitrarily small neighborhood of the shock front.

³For n = 1, g should be replaced by a constant, x_0 , and $M = \{x: x < x_0\}$.

THEOREM 3.12. Suppose d=n and u_0 is admissible and piecewise continuous. (See 3.5.) If $\varepsilon > 0$ there exists $T_{\varepsilon} > 0$ such that u(t,x) = v(x-kt) for $t \geq T_{\varepsilon}$ and $\operatorname{dist}(x-kt,\partial M) > \varepsilon$.

In §4 we study the case d < n and show that it reduces to the case described above. Without loss of generality we assume that $V \equiv \text{span}\{\tilde{f}(c): a \leq c \leq b\} = \mathbf{R}^d \times \{0\}$; the projection of μ on V is nontrivial; and its unit-length normalization, $(\mu', 0)$, satisfies

$$\langle (\mu', 0), \tilde{f}(c) \rangle \le -\theta(b-c)(c-a)$$
 for $a \le c \le b$.

(This holds after an appropriate change of coordinates.)

The problem, (0.4) and (0.5), can then be interpreted (after a finite time) as a parametrized set of problems in d space dimensions. To describe this in more detail, let us denote (x_1, \ldots, x_n) by (x', x'') with $x' \in \mathbf{R}^d$ and $x'' \in \mathbf{R}^{n-d}$. We prove

LEMMA 4.1. Suppose u_0 is admissible and d < n. Under the above hypotheses on \tilde{f} , there exists T > 0 (explicitly computable from \tilde{f}) so that if $t \geq T$,

$$u(t,x) = \tilde{u}(t,x'-k't;x''-k''t).$$

Here $\tilde{u}(t, x'; x'')$ is defined for a.e. $x'' \in \mathbf{R}^{n-d}$ as the solution of the problem

$$\frac{\partial \tilde{u}}{\partial t} + \sum_{i=1}^{d} \frac{\partial}{\partial x_i} (\tilde{f}_i(\tilde{u})) = 0 \quad \text{for } t \ge T; \ x' \in \mathbf{R}^d,$$
$$\tilde{u}(T, x'; x'') = u(T, x' + k'T, x'' + k''T).$$

The function u(T, x' + k'T, x'' + k''T) for fixed $x'' \in \mathbb{R}^{n-d}$ is admissible relative to a nondegenerate shock defined through μ' .

From our results for the case d = n, we have

THEOREM 4.6. Suppose u_0 is admissible, piecewise continuous, and d < n. For all $\varepsilon > 0$ and a.e. $z'' \in \mathbb{R}^{n-d}$ there exists $T = T(\varepsilon, z'') > 0$ and a Lipschitz domain, M(z''), in \mathbb{R}^d such that $\chi_M(z', z'') = \chi_{M(z'')}(z')$ and if x'' - k''t = z'',

$$u(t, x', x'') = v(t, x' - k't, x'' - k''t)$$

= $b \cdot \chi_{M(z'')}(x' - k't) + a \cdot \chi_{\mathbf{R}^d - M(z'')}(x' - k't)$

for $t \geq T$ and $dist(x' - k't, \partial M(z'')) \geq \varepsilon$.

Solutions of (0.4) and (0.5) when n=2 and u_0 is constant in each of the four quadrants were studied by D. Wagner in [9]. He assumed that $f=(f_1,f_2)$ with f_1'' , $f_2''>0$ and f_1 sufficiently close to f_2 . Under these assumptions he constructed the solution and analyzed its qualitative properties. Further study of this problem was done by W. B. Lindquist [6] assuming that $f_1=f_2$ and f_1 has at most two inflection points.

In [1] E. Conway studied the large-time behavior of solutions of (0.4) and (0.5) in n space dimensions assuming that u_0 has compact support and f is convex in some direction. He proved decay estimates on the L^{∞} -norm of such solutions.

As far as we know, all previous studies of the large-time behavior of u assumed some type of convexity condition on f. Our primary assumption is (0.6). In particular, our results are new even for the case n = 1.

1. Super and subentropy solutions. In this section we define a notion of super and subentropy solutions of a conservation law. A technique used throughout the paper is the explicit construction of such functions to obtain estimates on u(t, x) for large time.

First, let us recall some results proved by Vol'pert in [8]. Suppose $u \in BV(\mathbf{R}^{n+1}_+)$ $\cap L^{\infty}(\mathbf{R}^{n+1}_+)$. Then Vol'pert showed that with the exception of a set of H^n -measure zero every (t,x) in \mathbf{R}^{n+1} is a regular point of u. This means that there is a unit vector, ν , in \mathbf{R}^{n+1} and numbers $l_{\nu}u(t,x)$ and $l_{-\nu}u(t,x)$ such that for each $\varepsilon > 0$

$$\lim_{r\to 0}\frac{|\{(\tau,y)\in B_r(t,x)\colon |u(\tau,y)-l_{\pm\nu}u(t,x)|>\varepsilon \text{ and } \langle (\tau-t,y-x),\pm\nu\rangle\geq 0\}|}{|B_r(t,x)|}=0.$$

(Here |E| denotes the Lebesgue measure of a set, E.) We denote the set of regular points where $l_{\nu}u(t,x) \neq l_{-\nu}u(t,x)$ as the set of approximate jump points, $\Gamma(u)$. We remark that up to a multiple of ± 1 , ν is unique for $(t,x) \in \Gamma(u)$ and is called a normal to $\Gamma(u)$ at (t,x). We adopt the convention that $u^+ = l_{\nu}u(t,x)$, $u^- = l_{-\nu}u(t,x)$, and ν is chosen so that $u_+ \leq u_-$. If g(t) is a continuously differentiable function, Vol'pert proved that $g(u(t,x)) \in BV(\mathbf{R}_+^{n+1})$ and $\nabla g(u) = \tilde{g}(u) \cdot \nabla u$ as a Borel measure, where for each regular point, (t,x),

$$\tilde{g}(u) = \begin{cases} \frac{g(u^+) - g(u^-)}{u^+ - u^-} & \text{if } (t, x) \in \Gamma(u), \\ g'(\bar{u}(t, x)) & \text{if } (t, x) \not\in \Gamma(u), \end{cases}$$

and $\bar{u}(t,x) = \lim_{r\to 0} \int_{B_{\tau}(t,x)} u(\tau,y) d\tau dy$, the Lebesgue limit of u at (t,x). (This is well defined H^n -almost everywhere.)

Vol'pert proved that solutions of (0.4) and (0.5) are unique among those which satisfy an entropy condition. More precisely, suppose $u_0 \in BV(\mathbf{R}^n) \cap L^{\infty}(\mathbf{R}^n)$. Let A(u) = (u, f(u)) and define $S: \mathbf{R}^2 \to \mathbf{R}^{n+1}$ by

$$S(u, v) = (A(u) - A(v)) \cdot \operatorname{sign}(u - v).$$

Then there is a unique function $u \in BV(\mathbf{R}^{n+1}_+)$ satisfying

(1.1)
$$\partial u/\partial t + \operatorname{div} f(u) = \operatorname{div} A(u) = 0,$$

(1.2)
$$\lim_{t\to 0+} \bar{u}(t,x) = u_0(x) \quad \text{for } L^n\text{-a.e. } x$$

and the entropy condition:

(E)
$$\operatorname{div} S(u,c) \leq 0$$
 for all $c \in \mathbf{R}$ as a distribution on \mathbf{R}^{n+1}_+ .

When $u \in BV(\mathbf{R}^{n+1}_+) \cap L^{\infty}(\mathbf{R}^{n+1}_+)$ and u satisfies (1.1), the entropy condition, (E), is equivalent to

(E') $\langle S(u^+,c),\nu\rangle \leq \langle S(u^-,c),\nu\rangle$ for all $c\in \mathbf{R}$ and H^n -almost every point in the jump set, $\Gamma(u)$.

In addition, Vol'pert proved that if $u_0 \in BV(\mathbf{R}^n) \cap L^{\infty}(\mathbf{R}^n)$, then

- (i) $\bar{u}(t,x) \in C([0,\infty); L^1_{loc}(\mathbf{R}^n)),$
- $(ii) |\bar{u}| \leq ||u_0||_{\infty},$
- (iii) For any two solutions, u and w, of (1.1) which satisfy (E) and such that $u, w \in L^{\infty}(\mathbf{R}^{n+1}_+) \cap C([0,\infty); L^1_{\text{loc}}(\mathbf{R}^n))$ and $|u|, |w| \leq m < \infty$, we have

(1.3)
$$\int_{B_{\tau}} |\bar{u}(t,x) - \bar{w}(t,x)| dx \le \int_{B_{\tau+lt}} |u(x,0) - w(x,0)| dx$$

where $l = \left(\sum_{i=1}^n (\max_{|u| \leq m} |f_i'(u)|)^2\right)^{1/2}$.

If $u_0 \in L^{\infty}(\mathbf{R}^n)$, Vol'pert showed that there is a distribution solution of (1.1) and (1.2) in the class $C([0,\infty); L^1_{loc}(\mathbf{R}^n))$ which satisfies (E) and also satisfies (ii) and (iii). Kruzkov proved that it is unique in a broader class [5].

As stated in the introduction, we define a *solution* of (1.1) to be a function, u, in $C([0,\infty); L^1_{loc}(\mathbf{R}^n))$ which satisfies (1.1) in the sense of distributions and also satisfies the entropy condition (E).

We define super and subentropy solutions as follows:

DEFINITION 1.4. Suppose

$$u(t,x) \in BV(\mathbf{R}^{n+1}_+) \cap L^{\infty}(\mathbf{R}^{n+1}_+) \cap C([0,\infty); L^1_{\mathrm{loc}}(\mathbf{R}^n)).$$

Then u is a super (sub) solution of (1.1) if

(1)
$$\operatorname{div} A(u) \equiv \langle [A(u^{+}) - A(u^{-})], \nu \rangle d\mathcal{H}^{n}|_{\Gamma(u)} + \left(u_{t} + \sum_{i=1}^{n} f'_{i}(u) \cdot u_{x_{i}} \right) \Big|_{\Gamma(u)^{c}}$$
$$\geq 0 \ (< 0).$$

(2) For
$$\mathbb{X}^n$$
-a.e. $(t, x) \in \Gamma(u)$,
$$\langle A(u^+), \nu \rangle \geq (A(c), \nu) \quad \text{for } u^+ \leq c \leq u^-$$
$$(\text{sub: } \langle A(u^-), \nu \rangle \geq \langle A(c), \nu \rangle \text{ for } u^+ \leq c \leq u^-).$$

The inequalities in condition (2) are splitting of the entropy condition into two inequalities.

The following lemma will be used to prove a maximum principle for super and subsolutions. We use the notation $\sigma^+(z) = \chi_{\mathbf{R}^+}(z)$, $\sigma^-(z) = -\chi_{\mathbf{R}^-}(z)$, and $S^{\pm}(u,v) = [A(u) - A(v)] \cdot \sigma^{\pm}(u-v)$.

LEMMA 1.5. Suppose $u \in BV(\mathbf{R}^{n+1}_+) \cap L^{\infty}(\mathbf{R}^{n+1}_+) \cap C([0,\infty); L^1_{\mathrm{loc}}(\mathbf{R}^n))$. Then u is a supersolution iff $\operatorname{div} S^-(u,c) \leq 0$ for all $c \in \mathbf{R}$; u is a subsolution iff $\operatorname{div} S^+(u,c) \leq 0$ for all $c \in \mathbf{R}$.

PROOF. Let $\Delta S^-(u,c) = [S^-(u^+,c) - S^-(u^-,c)]$ in $\Gamma(u)$. Then from the calculus of BV functions [8, §15] we have

$$\operatorname{div} S^{-}(u,c) = \overline{\sigma^{-}}(u-c) \cdot \operatorname{div} A(u)|_{\Gamma(u)^{c}} + \langle \Delta S^{-}(u,c), \nu \rangle d\mathcal{H}^{n}|_{\Gamma(u)}.$$

Note first that $\langle \Delta S^-(u,c), \nu \rangle \leq 0$ a.e. on $\Gamma(u)$ for $c \in \mathbf{R}$ iff $\langle A(c), \underline{\nu} \rangle \leq \langle A(u^+), \nu \rangle$ for $u^+ \leq c \leq u^-$ a.e. on $\Gamma(u)$. Second, if we choose $c > ||u||_{\infty}$, then $\overline{\sigma^-}(u-c) \equiv -1$. Thus

$$\overline{\sigma^-}(u-c)\cdot \operatorname{div} A(u)|_{\Gamma(u)^c}\leq 0$$
 for all $c\in\mathbf{R}$

is equivalent to div $A(u) \ge 0$ in $\Gamma(u)^c$. This shows that div $S^-(u,c) \le 0$ iff u is a supersolution. The result for subsolutions is proved similarly.

THEOREM 1.6. Suppose u is a supersolution and v is a subsolution. Let

$$m=\sup\{\max\{|u(t,x)|,|v(t,x)|\}\colon x\in\mathbf{R}^n, 0\leq t\leq T\}$$

and

$$l = \left(\sum_{i=1}^n \left(\max_{|u| \le m} |f_i'(u)|\right)^2\right)^{1/2}.$$

Then div $S^-(u,v) = \text{div } S^+(v,u) \leq 0$ and for $0 \leq t \leq T < \infty$,

$$\int_{B_r} \max\{0, v - u\}(t, x) \, dx \le \int_{B_{r+lt}} \max\{0, v - u\}(0, x) \, dx.$$

PROOF. We have from [8, §15] that

$$\operatorname{div} S^{-}(u,v) = \overline{\sigma^{-}}(u-v) \cdot \operatorname{div}[A(u) - A(v)]|_{\Gamma(u-v)^{c}} + \langle \Delta S^{-}(u,v), \nu \rangle d\mathcal{H}^{n}|_{\Gamma(u-v)}$$

where $\Delta S^-(u,v) = S^-(l_{\nu}u,l_{\nu}v) - S^-(l_{-\nu}u,l_{-\nu}v)$ and ν is the normal to $\Gamma(u-v)$. The first term is a nonpositive measure. As for the second term, it follows from Lemma 1.5 that for any c

(1)
$$\langle S^-(l_{\nu}u,c),\nu\rangle \leq \langle S^-(l_{-\nu}u,c),\nu\rangle$$
 and

(2)
$$\langle S^+(l_{\nu}v,c),\nu\rangle \leq \langle S^+(l_{-\nu}v,c),\nu\rangle$$
,

 \mathcal{H}^n -almost everywhere on $\Gamma(u-v)$. Set $c=l_{\nu}v$ in (1) and $c=l_{-\nu}u$ in (2). Since $S^+(h,g)=S^-(g,h)$, it follows from (1) and (2) that $\Delta S^-(u,v)\leq 0$ a.e. on $\Gamma(u-v)$. This proves that div $S^-(u,v)=\operatorname{div} S^+(v,u)\leq 0$.

The second assertion follows from Green's theorem on the region, $D = \{(x, \tau): |x| < r + l(t - \tau), 0 < \tau < t\}$. We have (from [8, §14])

$$\begin{split} 0 &\geq \int_{D} \operatorname{div} S^{-}(u,v) \, d\tau \, dx = \int_{B_{\tau}} \max\{0,v-u\}(t,x) \, dx \\ &- \int_{B_{\tau+lt}} \max\{0,v-u\}(0,x) \, dx \\ &+ \int_{\partial D \cap \{0 < \tau < t\}} \langle S^{-}(l(u),l(v)),N \rangle d\mathcal{H}^{n} \end{split}$$

where l(u), l(v) are the inward traces of u, v and N is the outward-pointing normal on ∂D . From the choice of l the third term is nonnegative. This proves the theorem.

REMARK 1.7. The above theorem also holds if $v \notin BV(\mathbf{R}^{n+1}_+)$ but $v \in L^{\infty}(\mathbf{R}^{n+1}_+) \cap C([0,\infty); L^1_{loc}(\mathbf{R}^n))$ and v is a solution of (1.1). This follows from Vol'pert's construction of v as a limit in $C([0,\infty); L^1_{loc}(\mathbf{R}^n))$ of a sequence of functions, $\{v_j(t,x)\}$, which are solutions of (1.1) and satisfy the hypotheses of Theorem 1.6.

2. Preliminary results. In this section we prove preliminary results concerning the large-time behavior of solutions of (1.1) and (1.2). We assume throughout the paper that $f \in C^1(\mathbf{R}; \mathbf{R}^n)$ for $n \geq 1$, u(t,x) is a solution of (1.1) and (1.2), and u_0 is a compact perturbation of the Riemann data,

$$arphi(x) = \left\{ egin{aligned} a & ext{for } \langle x, \mu
angle > 0, \ b & ext{for } \langle x, \mu
angle < 0, \end{aligned}
ight.$$

where μ is a unit vector in \mathbf{R}^n and a < b.

We define

$$k = [f(b) - f(a)]/(b-a)$$
 and $\tilde{f}(c) = f(c) - f(a) - (c-a)k$.

Our main assumption (involving f and φ) is that φ is nondegenerate; that is, there exists a constant $\theta > 0$ such that

(2.1)
$$\langle \tilde{f}(c), \mu \rangle \leq -\theta(b-c)(c-a) \text{ for } a \leq c \leq b.$$

This is essentially a strict entropy condition on $\varphi(x-kt)$. To check this, we note that $w(t,x) \equiv \varphi(x-kt)$ is in $BV(\mathbf{R}^{n+1}_+) \cap L^{\infty}(\mathbf{R}^{n+1}_+)$ and is a weak solution of (1.1). The entropy condition (E) on w is equivalent to (E') which can be stated in the form

$$\langle A(c) - A(w^+), \nu \rangle \le \langle A(w^-) - A(c), \nu \rangle$$
 for $w^+ \le c \le w^-$

 H^n -a.e. in $\Gamma(w)$. (See §1.) Since ν is in the same direction as $(\langle -k, \mu \rangle, \mu)$ and $\tilde{f}(a) = \tilde{f}(b) = 0$, (E") reduces to

$$\langle \tilde{f}(c), \mu \rangle \leq 0$$
 for $a \leq c \leq b$.

Thus (2.1) ensures that $\varphi(x-kt)$ is a traveling wave solution of (1.1).

We define

$$m_1 = \inf \left\{ m \colon m < a ext{ and } lpha_1(m) \equiv \max_{m \leq c \leq a} rac{\langle ilde{f}(c), \mu
angle}{c - a} < eta_1(m) \equiv \min_{m \leq c \leq a} rac{\langle ilde{f}(c), \mu
angle}{c - b}
ight\}$$

and

$$m_2 = \sup \left\{ m \colon m > b \text{ and } \alpha_2(m) \equiv \max_{b \leq c \leq m} \frac{\langle \tilde{f}(c), \mu \rangle}{c - a} < \beta_2(m) \equiv \min_{b \leq c \leq m} \frac{\langle \tilde{f}(c), \mu \rangle}{c - b} \right\}.$$

By (2.1) and the differentiability of f, m_1 and m_2 are well defined with $-\infty \le m_1 < a$ and $b < m_2 \le +\infty$. The initial data, u_0 , is called *admissible* if

- (i) for some $R_1 > 0$, $u_0(x) = \varphi(x)$ for $|x| \ge R_1$, and
- (ii) $u_0 \in L^{\infty}(\mathbf{R}^n)$, $m_1 < \text{ess inf } u_0$, and $m_2 > \text{ess sup } u_0$.

We first prove that if u_0 is admissible then for some $t^* < \infty$, we have $a \le u(t,x) \le b$ for $t \ge t^*$. It will be convenient in the proof to assume (without loss of generality) that f(a) = f(b) = 0 and hence k = 0. This is possible because if $k \ne 0$ we can make a change of variables,

$$(2.2) x_1 = x - kt, t_1 = t,$$

and set $\tilde{u}(t_1, x_1) = u(t, x)$. Equations (1.1), (1.2), and the entropy condition (E) become

$$\partial \tilde{u}/\partial t_1 + \operatorname{div}_{x_1} \tilde{f}(\tilde{u}) = 0, \quad \tilde{u}(0,x) = u_0(x),$$

and

(
$$\tilde{\mathbf{E}}$$
) $\operatorname{div}_{(t_1,x_1)}\tilde{S}(\tilde{u},c) \leq 0$ for all $c \in \mathbf{R}$

where $\tilde{S}(v,w) = \text{sign}(v-w) \cdot [\tilde{A}(v) - \tilde{A}(w)]$ and $\tilde{A}(v) = (v, \tilde{f}(v))$. Thus the effect is to replace f by \tilde{f} which satisfies $\tilde{f}(a) = \tilde{f}(b) = 0$.

THEOREM 2.3. If u_0 is admissible, there exists $t^* < \infty$ (defined below) such that $a \le u(t, x) \le b$ for $t \ge t^*$.

PROOF. Let $b' = \operatorname{ess\ sup} u_0$ and $a' = \operatorname{ess\ inf} u_0$. Define

$$\begin{split} t_1 &= \left\{ \begin{aligned} &0 & \text{if } a' \geq a, \\ &2R_1/[\beta_1(a') - \alpha_1(a')] & \text{if } a' < a, \\ &t_2 &= \left\{ \begin{aligned} &0 & \text{if } b' \leq b, \\ &2R_1/[\beta_2(b') - \alpha_2(b')] & \text{if } b' > b, \end{aligned} \right. \end{split}$$

and let $t^* = \max\{t_1, t_2\}.$

First we show that $u \leq b$ for $t \geq t^*$. We assume that b' > b because if not, $u \leq b$ for all time by the maximum principle. We also assume (w.l.o.g.) that f(a) = f(b) = 0 and hence $f = \tilde{f}$.

Let $\alpha = \alpha_2(b')$, $\beta = \beta_2(b')$, and consider the function defined by

$$w(t,x) = \begin{cases} b & \text{for } \langle x,\mu \rangle < \beta t - R_1, \\ b' & \text{for } \beta t - R_1 < \langle x,\mu \rangle < \alpha t + R_1, \\ a & \text{for } \alpha t + R_1 < \langle x,\mu \rangle, \end{cases}$$

when $0 \le t \le 2R_1/(\beta - \alpha)$;

$$w(t,x) = \left\{ egin{array}{ll} b & ext{for } \langle x,\mu
angle < 2R_1lpha/(eta-lpha) + R_1, \ a & ext{for } 2R_1lpha/(eta-lpha) + R_1 < \langle x,\mu
angle, \end{array}
ight.$$

when $t > 2R_1/(\beta - \alpha)$.

We claim that w is a supersolution. Clearly $w \in BV(\mathbf{R}^{n+1}_+) \cap L^{\infty}(\mathbf{R}^{n+1}_+) \cap C([0,\infty); L^1_{loc}(\mathbf{R}^n))$ and since w is piecewise constant we need only verify (1.4.2) on $\Gamma(w)$.

On the jump from w=b to w=b' for $t\leq 2R_1/(\beta-\alpha)$, we have $\nu=(\beta,-\mu)/\sqrt{1+\beta^2}$. We need

$$\langle A(b) - A(c), \nu \rangle \ge 0$$
 for $b \le c \le b'$.

The left-hand side has the same sign as $-(c-b)\beta + \langle f(c), \mu \rangle$ and this is nonnegative for $b \leq c \leq b'$ by the definition of β . On the jump from w = b' to w = a, we take ν in the direction of $(-\alpha, \mu)$. It is required that $\langle A(a) - A(c), \nu \rangle \geq 0$ for $a \leq c \leq b'$. This is equivalent to

$$\alpha(c-a) - \langle f(c), \mu \rangle \ge 0$$
 for $a \le c \le b'$

which holds by (2.1) and the definition of α since $\alpha \geq \langle f(b), \mu \rangle / (b-a) = 0$.

The entropy condition holds across the jump from w=b to w=a for $t>2R_1/(\beta-\alpha)$ since $\nu=(0,\mu)$ and thus the entropy condition (E') reduces to $\langle f(c),\mu\rangle\leq 0$ for $a\leq c\leq b$. By Remark 1.7, we have

$$\int_{B_r} \max\{0, w - u\}(t, x) \, dx \le \int_{B_{r+1t}} \max\{0, w - u\}(0, x) dx = 0$$

for any t and $r \geq 0$ where $l = \left[\sum_{i=1}^n (\max_{|u| \leq |b'| + |a'|} |f_i'(u)|)^2\right]^{1/2}$. We conclude that $u \leq w$ for all t and a.e. $x \in \mathbf{R}^n$, and hence $u \leq b$ for $t \geq t^*$ and a.e. $x \in \mathbf{R}^n$. To show that $u \geq a$ when $t \geq t^*$ involves the construction of a subsolution in a similar way.

A consequence of Theorem 2.3 is the following result on the large-time behavior of u.

THEOREM 2.4. Suppose u_0 is admissible. There exists $R_3 > 0$ and a set $M \subset \mathbb{R}^n$ such that if $v(x) = b \cdot \chi_M(x) + a \cdot \chi_{\mathbb{R}^n - M}(x)$, then

- (i) $u(t,x) v(x-kt) \rightarrow 0$ in $L^1(\mathbf{R}^n)$ as $t \rightarrow \infty$,
- (ii) u(t, x) = v(x kt) for t > 0 and $|x kt| > R_3$, and
- (iii) $M \cap \{x \in \mathbf{R}^n : |x| > R_3\} = \{x \in \mathbf{R}^n : \langle x, \mu \rangle < 0 \text{ and } |x| > R_3\}.$

To prove this, we will need the following lemma.

LEMMA 2.5. Suppose u_0 is admissible. There exists $R_2 > R_1$ such that $u(t, x) = \varphi(x - kt)$ if $t \ge 0$ and $|x - kt| \ge R_2$.

PROOF. By the change of variables (2.2), it is sufficient to prove the lemma assuming f(a) = f(b) = 0, so that k = 0 and $f = \tilde{f}$. Since φ is then a steady state solution of (1.1) and $u_0 = \varphi$ for $|x| \geq R_1$, it follows from estimate (1.3) that for some $r_0 > R_1$,

$$u(t^*, x) = \varphi(x)$$
 for $|x| \ge r_0$

where t^* is as defined in Theorem 2.3. We will show that in fact there exists $r_1 > r_0$ such that

$$u(t,x) = \varphi(x)$$
 for $t \ge t^*$, $|x| \ge r_1$.

To prove this we use the nondegeneracy condition, (2.1). It ensures that for some small $\delta > 0$, $\langle f(c), \nu \rangle \leq 0$ for $a \leq c \leq b$ and all $\nu \in \mathbf{R}^n$ such that $A(\mu, \nu) < \delta$, where $A(\mu, \nu)$ is the angle between μ and ν that is not greater than π . Hence we can construct open sets, M_1 and M_2 (with smooth boundaries) and choose $r_1 > r_0$ such that

$$\begin{split} &M_1\subset \{x\colon\!\langle x,\mu\rangle<0\}\subset M_2,\\ &B_{r_0}\equiv B_{r_0}(0)\subset M_2-M_1,\\ &M_i\cap \{x\colon\!|x|\geq r_1\}=\{x\colon\!\langle x,\mu\rangle<0\text{ and }|x|\geq r_1\}\quad\text{for }i=1,2,\text{ and }\\ &\langle f(c),\nu\rangle\leq 0\quad\text{for }a\leq c\leq b, \end{split}$$

where ν is any outward pointing normal to M_1 or M_2 . As a result the functions, $v_i(x) = b \cdot \chi_{M_i}(x) + a \cdot \chi_{\mathbf{R}^n - M_i}(x)$, are steady state solutions of (1.1). Since $v_1(x) \leq u(t^*, x) \leq v_2(x)$ for all x, we have

$$v_1(x) \le u(t,x) \le v_2(x)$$
 for $x \in \mathbb{R}^n$, $t \ge t^*$.

We conclude that

$$u(t,x) = \varphi(x)$$
 for $|x| \ge r_1$, $t \ge t^*$.

By (1.3) the above equation holds when $t \ge 0$ and $|x| \ge R_2 \equiv r_1 + lt^*$ which proves the conclusion of the lemma.

PROOF OF THEOREM 2.4. As before, we may assume that f(a)=f(b)=0 and hence k=0.

Suppose $u_0 \in BV(\mathbf{R}^n)$. Consider $u_h(t,x) = u(t,x+he_i)$ for any $i=1,\ldots,n$ and $|h| \leq R_1$, where $\{e_1,\ldots,e_n\}$ is the standard basis for \mathbf{R}^n . Since u_h is a solution of (1.1), we have from Vol'pert's results (or Theorem 1.6) that $\operatorname{div} S(u_h,u) \leq 0$. Let $R=3 \cdot R_2$ with R_2 as defined in Lemma 2.5 and $B_R=B_R(0)$. By applying Green's theorem in the cylinder, $(0,t) \times B_R$, we obtain

$$\begin{split} \int_{B_R} |u(t,x+he_i) - u(t,x)| dx &\leq \int_{B_R} |u(0,x+he_i) - u(0,x)| dx \\ &- \int_0^t \int_{\partial B_R} \left\langle \frac{x}{R}, ([f(u_h) - f(u)] \cdot \operatorname{sign}(u_h - u)) \right\rangle d\mathcal{X}^n. \end{split}$$

By Lemma 2.5, $f(u_h) = f(u) = 0$ on ∂B_R for any t. Hence the third term is zero and

$$\int_{B_R} \left| \frac{\partial u}{\partial x_i}(t, x) \right| dx \le \int_{B_R} \left| \frac{\partial u}{\partial x_i}(0, x) \right| dx.$$

This and equation (1.1) imply that

$$\int_{-1}^{1} \int_{B_{R}} |(Du)(s+t,x)| dx \, ds \le c \cdot \int_{B_{R}} |(\nabla_{x}u)(0,x)| dx$$

for any t > 1, where $D = (\partial/\partial t, \partial/\partial x_1, \dots, \partial/\partial x_n)$.

Now multiply equation (1.1) by $\langle x, \mu \rangle$ and integrate over the cylinder $(t^*, T) \times B_R$. We obtain

$$\begin{split} \int_{B_R} [u(T,x) - u(0,x)] \cdot \langle x, \mu \rangle dx \\ &= \int_{t^*}^T \int_{B_R} \langle f(u), \mu \rangle dx \, dt - \int_{\partial B_R} \left\langle f(u), \frac{x}{R} \right\rangle \cdot \langle x, \mu \rangle d\mathcal{X}^n. \end{split}$$

The first term is bounded uniformly in T (by estimate (1.3)) and the last term is zero. Since $a \le u \le b$ for $t \ge t^*$, we conclude using (2.1) that

$$\int_0^\infty \int_{B_R} |(u-a)(u-b)| dx \, dt < \infty.$$

By Lemma 2.5, this can be restated as

$$\int_{\mathbf{R}^{n+1}_{\perp}} |(u-a)(u-b)| dx dt < \infty.$$

From the above estimates (and Lemma 2.5) we see that for some sequence, $t_j \to +\infty$, $\{u(s+t_j,x)\}$ converges in $L^1_{loc}((-1,1)\times \mathbf{R}^n)$ to a function, $v(s,x)\in BV((-1,1)\times \mathbf{R}^n)$. The L^1 bound on (u-a)(u-b) implies that (v-a)(v-b)=0 pointwise almost everywhere in $(-1,1)\times \mathbf{R}^n$. Thus

$$v(s,x) = b \cdot \chi_{\mathbf{M}}(s,x) + a \cdot \chi_{\mathbf{R}^n - \mathbf{M}}(s,x)$$

with $\chi_M \in BV((-1,1) \times \mathbf{R}^n)$. On the other hand, div $S(v,c) \leq 0$ as a measure on $(-1,1) \times \mathbf{R}^n$ for all $c \in \mathbf{R}$. Hence v is a solution of (1.1) on $(-1,1) \times \mathbf{R}^n$. (See [8, §16].) Since f(a) = f(b) = 0, we have

$$0 = \partial v/\partial s + \operatorname{div} f(v) = \partial v/\partial s$$
 on $(-1, 1) \times \mathbf{R}^n$.

Thus $\chi_M(s,x) = \chi_M(x)$.

Finally we show that $u(t,x) - v(x) \to 0$ in $L^1(\mathbf{R}^n)$ as $t \to \infty$. By Lemma 2.5 it is sufficient to prove convergence in $L^1(B_R)$. This follows from estimate (1.3) which implies that

$$\int_{B_R(0)} |v(x)-u(t,x)| dx \leq \int_{B_R} |v(x)-u(s+t_j,x)| dx$$

for $t \geq s + t_j$. This proves (i) when $u_0 \in BV(\mathbf{R}^n)$; (ii) and (iii) follow from (i) and Lemma 2.5 with $R_3 = R = 3 \cdot R_2$.

For the case when $u_0 \notin BV(\mathbf{R}^n)$ we can approximate u_0 by admissible functions, u_j , in $BV(\mathbf{R}^n)$ such that $u_j \to u_0$ in $L^1(B_{R_3})$ and $u_j(x) = \varphi(x)$ for $|x| \ge R_1$. By Lemma 2.5 if $u_j(t,x)$ is the solution of (1.1) with initial values, $u_j(x)$, then

$$u_j(t,x) = u(t,x) = \varphi(x)$$
 for $|x| \ge R_2$.

By estimate (1.3) it follows that

$$\int_{B_{R_3}} |v_i(x) - v_j(x)| dx \le \int_{B_{R_3}} |u_i(x) - u_j(x)| dx$$

where $v_j(x) = \lim_{t\to\infty} u_j(t,x)$. Hence for some set M, $\chi_{M_j} - \chi_M \to 0$ in $L^1(\mathbf{R}^n)$ and so

$$u(t,x) - b \cdot \chi_{\mathbf{M}}(x) - a \cdot \chi_{\mathbf{R}^n - \mathbf{M}}(x) \to 0$$

in $L^1(\mathbf{R}^n)$ as $t \to \infty$.

3. The *n*-dimensional case. We now get a more precise description of the set M and the convergence of u to v(x-kt) in the special case

(3.1)
$$n = d \equiv \dim \operatorname{span}\{\tilde{f}(c) : a \le x \le b\}.$$

THEOREM 3.2. There is a system of coordinates, (\hat{y}, y_n) , obtained by a rotation and a constant C_0 (both depending only on \tilde{f}) so that if u_0 is admissible then $M = \{(\hat{y}, y_n): y_n < g(\hat{y})\}$. The shock front, g, is Lipschitz continuous with Lipschitz constant at most C_0 .

PROOF. First assume $u_0 \in BV(\mathbf{R}^n)$ and w.l.o.g. f(a) = f(b) = 0 so that $f = \tilde{f}$. From the proof of Theorem 2.4, v(x) is a steady state solution of (1.1) and $v \in BV(\mathbf{R}^n)$. Hence by $(E') \langle f(c), \nu(x) \rangle \leq 0$ for $a \leq c \leq b$ and \mathcal{A}^{n-1} a.e. $x \in \Gamma(v)$ where $\nu(x)$ is the normal to $\Gamma(v)$ at x. Consider the convex set

$$\Omega = \left\{ \sum_{i=1}^{m} a_i f(c_i) : a_i \ge 0, \ a \le c_i \le b, \ m < \infty \right\}.$$

We have $\nu(x) \in \bigcap_{z \in \Omega} \{w \in \mathbb{R}^n : \langle w, z \rangle \leq 0\}$ for each such $x \in \Gamma(v)$. From (3.1) the interior of Ω is nonempty. Choose $\alpha \in \Omega^0$ with $|\alpha| = 1$ and $\delta > 0$ so that

$$E \equiv \{z \in \mathbf{R}^n : A(z, \alpha) \le \delta\} \subset \Omega.$$

Then

$$\nu(x) \in \bigcap_{z \in E} \{ w \in \mathbf{R}^n : \langle w, z \rangle \le 0 \} = \{ w \in \mathbf{R}^n : A(-\alpha, w) \le \pi/2 - \delta \}.$$

It follows from the theory of sets with locally finite perimeter [3, Theorem 4.8] that in a rotated system of coordinates (with $e_n = -\alpha$) there exists a function $g(\hat{y})$ such that $M = \{(\hat{y}, y_n): y_n < g(\hat{y})\}$ where $v = b \cdot \chi_M + a \cdot \chi_{\mathbb{R}^n - M}$ and

$$|g(\hat{y}_1) - g(\hat{y}_2)| \le \tan(\pi/2 - \delta) \cdot |\hat{y}_1 - \hat{y}_2| \equiv C_0 \cdot |\hat{y}_1 - \hat{y}_2|.$$

If $u_0 \not\in BV(\mathbf{R}^n)$ we take a sequence, $\{u_j(0,x)\}$, of admissible data such that $u_j(0,x) \in BV(\mathbf{R}^n)$ and $u_j(0,x) \to u_0(x)$ in $L^1_{\mathrm{loc}}(\mathbf{R}^n)$ as $j \to \infty$. We have seen that the corresponding steady state limits satisfy $v_j \to v$ in $L^1_{\mathrm{loc}}(\mathbf{R}^n)$ as $j \to \infty$. From the above argument we have $M_j = \{(\hat{y},y_n): y_n < g_j(\hat{y})\}$ where the g_j are uniformly Lipschitz continuous and

$$\langle f(c), (-\nabla g_j(\hat{y}), 1) \rangle \leq 0 \text{ for } a \leq c \leq b$$

and L^{n-1} a.e. \hat{y} . (Here it is understood that f(c) is expressed in the new coordinates, (\hat{y}, y_n) .) It follows that $M = \{(\hat{y}, y_n) : y_n < g(\hat{y})\}$ for some Lipschitz continuous g satisfying the same conditions.

The uniform bound on the Lipschitz constant of the shock front yields the following stability result.

THEOREM 3.3. Let $u_1(0,x)$, $u_2(0,x)$ be admissible data and $v_1(x-kt)$, $v_2(x-kt)$ be the asymptotic limits with g_1 , g_2 the corresponding shock fronts. There is a constant $C_1(\tilde{f}) < \infty$ so that

$$||g_1-g_2||_{L^{\infty}(\mathbf{R}^{n-1})} \le C_1 \left(\int_{B_{R_1}} |u_1(0,x)-u_2(0,x)| dx \right)^{1/n}.$$

PROOF. Assume that f(a) = f(b) = 0. In the coordinates (\hat{y}, y_n) we have $M_i = \{(\hat{y}, y_n): y_n < g_i(\hat{y})\}$ and $|\nabla g_i| \le C_0$ for i = 1, 2 where $v_i = b \cdot \chi_{M_i} + a \cdot \chi_{\mathbf{R}^n - M_i}$. As in the proof of Theorem 2.4, Green's theorem implies that

$$(b-a) \cdot \int_{\mathbf{R}^{n-1}} |g_1 - g_2| d\hat{y} = \int_{B_{R_3}} |v_1 - v_2| dx$$

$$\leq \int_{B_{R_1}} |u_1(0, x) - u_2(0, x)| dx.$$

But since $|\nabla g_1|$ and $|\nabla g_2|$ are a priori bounded, we conclude that if $|g_1(\hat{y}_0) - g_2(\hat{y}_0)| = \delta > 0$ then $|g_1(\hat{y}) - g_2(\hat{y})| \ge \delta/2$ for $\hat{y} \in B_{\delta/4C_0}(\hat{y}_0)$. Thus

$$(||g_1-g_2||_{L^{\infty}(\mathbb{R}^{n-1})})^n \le C_1(\tilde{f}) \cdot \int_{\mathbb{R}^{n-1}} |g_1-g_2| d\hat{y}.$$

REMARK 3.4. We will need to apply a version of this theorem when $u_1(0,x)$ and $u_2(0,x)$ are admissible relative to shocks that differ by a small translation, i.e. $u_1(0,x)=\varphi(x)$ and $u_2(0,x)=\varphi(x+\eta\mu)$ for $|x|\geq R_1$. If we assume w.l.o.g. that $C_0>1$ and take $R_4=5C_0\cdot\max\{R_3,C_0\eta\}$ then the same argument yields

$$||g_1-g_2||_{L^{\infty}(\mathbf{R}^{n-1})} \le C_1 \left(\int_{B_{R_4}} |u_1(0,x)-u_2(0,x)| dx \right)^{1/n}.$$

DEFINITION 3.5. A function h(x) is piecewise continuous iff there exists a finite number of mutually disjoint domains, $\{D_i\}_{i=1}^k$, such that $h|_{D_i}$ has a continuous extension to $D_i \bigcup \partial D_i$ and $L^n\left(\mathbf{R}^n - \bigcup_{i=1}^k D_i\right) = 0$.

We will prove that if u_0 is admissible, piecewise continuous, and $\operatorname{dist}(x-kt,\partial M) > \varepsilon > 0$, then u(t,x) = v(x-kt) after a finite time depending on ε and \tilde{f} . (See Theorem 3.12.) It will be convenient to use the following notation.

DEFINITION 3.6. If

$$w(t,x) \in C([0,\infty); L^1_{\mathrm{loc}}(\mathbf{R}^n))$$

we denote by U(w)(t,x) (L(w)(t,x)) the upper (lower) Lebesgue limit in x:

$$U(w)(t,x) = \overline{\lim}_{r \to 0} \int_{B_r(x)} w(t,z)dz, \quad L(w)(t,x) = \underline{\lim}_{r \to 0} \int_{B_r(x)} w(t,z)dz.$$

Note that for each t > 0, U(w)(t, x) = L(w)(t, x) = w(t, x) for a.e. x.

The following lemma implies the uniform convergence of u(t,x) away from the shock front.

LEMMA 3.7. Suppose that u_0 is admissible, piecewise continuous, and $g(\hat{y})$ is the shock front for the asymptotic limit. Then given ε , $\sigma > 0$ there is a constant $T(\varepsilon, \sigma) < \infty$ so that in the coordinates (\hat{y}, y_n) for t > T we have

- (1) $a \le u(t, y) \le a + \sigma$ for a.e. y with $y_n k_n t \ge g(\hat{y} \hat{k}t) + \varepsilon$,
- (2) $b \ge u(t,y) \ge b \sigma$ for a.e. y with $y_n k_n t \le g(\hat{y} \hat{k}t) \varepsilon$, where (\hat{k}, k_n) is the representation of k = [f(b) f(a)]/(b-a) in the coordinates, (\hat{y}, y_n) .

PROOF. Assume f(a) = f(b) = 0 so that k = 0. Note that from Lemma 2.3 the first inequalities in (1) and (2) hold for $t \ge t^*$.

For any $\delta \in (0,1)$ we can find functions $q_{\delta}(x)$, $p_{\delta}(x)$ and a positive constant, $\tau(\delta) \leq \delta$, so that the following three properties hold. First,

ess inf
$$u_0 \le q_{\delta}(x) \le u_0(x+h) \le p_{\delta}(x) \le \text{ess sup } u_0$$

for any $h \in \mathbf{R}^n$ with $|h| \leq \tau$. Second, if $q_{\delta}(t,x)$ and $p_{\delta}(t,x)$ are the solutions of (1.1) with initial values $p_{\delta}(x)$ and $q_{\delta}(x)$, we have $q_{\delta}(t,x) = \varphi(x+\tau\mu)$ and $p_{\delta}(t,x) = \varphi(x-\tau\mu)$ for $|x| \geq R$, $t \geq 0$, where R is independent of δ . Third,

$$\int_{B_{R}}(p_{\delta}-q_{\delta})dx\leq\delta.$$

If (1) does not hold there exists a sequence of points (t^m, y^m) in the rotated coordinates, $y = (\hat{y}, y_n)$, with $t^m \to \infty$ and $y^m \to y^0$ as $m \to \infty$ so that

$$U(u)(t^m,y^m)>a+\sigma\quad\text{and}\quad y_n^m\geq g(\hat{y}^m)+\varepsilon.$$

For δ fixed and any $\bar{y} \in B_{\tau/2}(y^0)$ it follows from Remark 1.7 that

$$u(t^m, y + h) \le p_{\delta}(t^m, y)$$
 for a.e. y

where $h = y^m - \bar{y}$, m is sufficiently large and $p_{\delta}(t, y)$ denotes the function $p_{\delta}(t, x)$ expressed in the rotated spacial coordinates. Thus

$$U(u)(t^m, y^m) \leq U(p_{\delta})(t^m, \bar{y}).$$

As a result,

(3)
$$a + \sigma \le p_{\delta}(t^m, \bar{y})$$
 for a.e. $\bar{y} \in B_{\tau/2}(y^0)$

if m is large. Let $v_{\delta}(x)$ be the asymptotic limit of $p_{\delta}(t,x)$,

$$v_{\delta}(x) = b \cdot \chi_{M_{\delta}}(x) + a \cdot \chi_{\mathbf{R}^n - M_{\delta}}(x).$$

It follows from (3) and Theorem 2.4 that $B_{\tau/2}(y^0) \subset M_{\delta}$. Using $g_{\delta}(\hat{y})$ to denote the shock front for v_{δ} we get $\varepsilon \leq |g_{\delta}(\hat{y}^0) - g(\hat{y}^0)|$. But from (3.4)

$$||g-g_{\delta}||_{L^{\infty}(\mathbf{R}^{n-1})} \leq C \left(\int_{B_{R_4}} |u_0(x)-p_{\delta}(0,x)| dx \right)^{1/n} \leq C \cdot \delta^{1/n}.$$

Since δ can be chosen arbitrarily small this is a contradiction. The verification of (2) is a similar argument using $q_{\delta}(x)$.

We shall improve on this estimate by using super and subsolutions. The construction of the supersolution is done below in detail.

LEMMA 3.8. Let f(c) be expressed in the coordinates (\hat{y}, y_n) and suppose f(a) = f(b) = 0. Let $g(\hat{y})$ be any Lipschitz continuous function so that for almost every \hat{y} ,

$$(3.9) \qquad \frac{\langle (\nabla g, -1), f(c) \rangle}{c - a} \ge m > 0 \quad \text{for } a \le c \le \frac{a + b}{2}, \\ \langle (\nabla g, -1), f(c) \rangle \ge 0 \qquad \quad \text{for } a \le c \le b.$$

Suppose η and δ are positive constants. Define

$$w(t, \hat{y}, y_n) = \begin{cases} b & \text{for } y_n \leq g(\hat{y}) - \delta + \eta t, \\ (a+b)/2 & \text{for } g(\hat{y}) - \delta + \eta t \leq y_n \leq g(\hat{y}) - mt, \\ a & \text{for } g(\hat{y}) - mt \leq y_n, \end{cases}$$

when $0 \le t \le \delta/(m+\eta)$;

$$w(t, \hat{y}, y_n) = \begin{cases} b & \text{for } y_n < g(\hat{y}) - m\delta/(m + \eta), \\ a & \text{for } y_n \ge g(\hat{y}) - m\delta/(m + \eta) \end{cases}$$

when $t \geq \delta/(m+\eta)$.

Then w is a supersolution of (1.1) provided η is sufficiently large depending only on f and the Lipschitz constant of q.

PROOF. Since w is piecewise constant,

$$\partial_t w + \operatorname{div}_x f(w) = \langle |[(w^+, f(w^+)) - (w^-, f(w^-))], \nu(t, y) \rangle d\mathcal{H}^n|_{\Gamma(w)}.$$

We must verify that

$$\langle [(w^+, f(w^+)) - (c, f(c))], \nu \rangle \geq 0$$
 for $w^+ \leq c \leq w^-$

and \mathcal{H}^n a.e. $(t,y) \in \Gamma(w)$. The jump between the states b and (a+b)/2 is determined by

$$0 = y_n - g(\hat{y}) + \delta - \eta t.$$

We have

$$\nu = (-\eta, -\nabla g, 1) / \sqrt{\eta^2 + |\nabla g|^2 + 1}.$$

Thus we need to establish that

$$-\left(\frac{a+b}{2}-c\right)\eta+\left\langle \left[f\left(\frac{a+b}{2}\right)-f(c)\right],(-\nabla g,1)\right\rangle \geq 0\quad\text{for }\frac{a+b}{2}\leq c\leq b.$$

This is equivalent to

$$\eta \geq \left\langle \left[f\left(rac{a+b}{2}
ight) - f(c)
ight], (-
abla g, 1)
ight
angle \left/ \left(rac{a+b}{2} - c
ight)$$

which is valid for η sufficiently large. For the jump between (a+b)/2 and a we have $0=y_n-g(\hat{y})+mt,\ \nu=(m,-\nabla g,1)/\sqrt{m^2+|\nabla g|^2+1}$. We must verify that

$$m(a-c) - \langle f(c), (-\nabla g, 1) \rangle \ge 0$$
 for $a \le c \le (a+b)/2$.

This is equivalent to the first inequality of (3.9).

On the jump between b and a for $t \ge \delta/(m+\eta)$ the entropy condition is equivalent to the second inequality of (3.9). This completes the proof.

The shock front q satisfies

$$g(\hat{y}) = \langle l, \hat{y} \rangle$$
 for $|\hat{y}| \geq R_3$

where $(l,-1)/\sqrt{|l|^2+1}$ is the representation of $-\mu$ in the coordinates (\hat{y},y_n) . Hence

$$\langle f(c), (l, -1) \rangle \ge \theta(c - a)(b - c)\sqrt{|l|^2 + 1}$$
 for $a \le c \le b$.

In general the first inequality of (3.9) does not hold. However, g can be uniformly approximated from above and below by functions, $g_{1,\varepsilon}$ and $g_{2,\varepsilon}$ respectively, where

$$g_{j,\varepsilon}(\hat{y}) = g((1-\varepsilon)\hat{y}) + \varepsilon \langle l, \hat{y} \rangle - (-1)^j \cdot c\varepsilon$$

for $0 < \varepsilon < \frac{1}{2}, \ j = 1, 2$, and c > 0 is sufficiently large (independent of ε). We have

$$\nabla g_{j,\varepsilon}(\hat{y}) = (1-\varepsilon) \cdot \nabla g((1-\varepsilon)\hat{y}) + \varepsilon l$$

so that

(3.10)
$$\langle f(c), (\nabla g_{j,\varepsilon}(\hat{y}), -1) \rangle$$

$$= (1 - \varepsilon) \langle f(c), (\nabla g((1 - \varepsilon)\hat{y}), -1) \rangle + \varepsilon \langle f(c), (l, -1) \rangle$$

$$\geq \varepsilon \theta \sqrt{|l|^2 + 1} \cdot (c - a)(b - c) \quad \text{for } a \leq c \leq b.$$

Thus supersolutions can be constructed with $g_{1,\varepsilon}$.

Similarly one can construct subsolutions with $g_{2,\varepsilon}$. The inequalities analogous to (3.9) which we require are

$$\frac{\langle (\nabla g, -1), f(c) \rangle}{b - c} \ge \eta > 0 \quad \text{for } \frac{a + b}{2} \le c \le b,$$
$$\langle (\nabla g, -1), f(c) \rangle \ge 0 \quad \text{for } a \le c \le b.$$

These hold with g replaced by $g_{2,\varepsilon}$ and η sufficiently small by inequality (3.10). Thus the formula of Lemma 3.8 defines a subsolution if m is sufficiently large.

We now improve on Lemma 3.7 in the sense that we show u(t,x) = v(x - kt) outside of an arbitrarily small neighborhood of the shock front after finite time.

THEOREM 3.11. Let u(t,x) be a solution of (1.1) with $u_0(x)$ admissible and piecewise continuous. Then in the coordinates (\hat{y}, y_n) , given $\delta > 0$ there is a constant $T(\delta) < \infty$ so that

(1)
$$u(t,y) = a \text{ for a.e. } y \text{ with } y_n - k_n t \ge g(\hat{y} - \hat{k}t) + \delta, \ t \ge T,$$

(2)
$$u(t,y) = b \text{ for a.e. } y \text{ with } y_n - k_n t \le g(\hat{y} - \hat{k}t) - \delta, \ t \ge T.$$

PROOF. Assume that f(a) = f(b) = 0. Suppose there is a sequence (t^j, y^j) with $y^j \to y^0$, $t^j \to \infty$ as $j \to \infty$ such that $U(u)(t^j, y^j) > a$ and $y_n^0 \ge g(\hat{y}^0) + \delta$. Suppose $\varepsilon \ll \delta$ and set

$$h_{\tau}(\hat{y}) = g_{1,\varepsilon}(\hat{y}) + \tau, \quad M_{\tau} = \{(\hat{y}, y_n): y_n < h_{\tau}(\hat{y})\}.$$

Define

$$\tau(t) = \inf\{\tau : \{y : U(u)(t,y) > a\} \subset M_{\tau}\}.$$

This is well defined for each t > 0 by Theorem 2.4. Set $v_{\tau} = b \cdot \chi_{M_{\tau}} + a \cdot \chi_{\mathbf{R}^n - M_{\tau}}$. By (3.10), v_{τ} is a steady state solution of (1.1). From Remark 1.7 and Theorem 2.3 we get $v_{\tau(t)}(y) \geq u(s,y)$ for $s \geq t$ and a.e. y. Thus $\tau(t)$ is nonincreasing as $t \uparrow \infty$.

Set $\tau_0 = \lim_{t\to\infty} \tau(t)$. If ε is sufficiently small $y_n^0 \ge h_{3\delta/4}(\hat{y}^0)$, so $\tau_0 \ge 3\delta/4$. Hence $h_{\tau_0} \ge g + 3\delta/4$. Choose T so large that $\tau(T) \le \varepsilon^2 + \tau_0$; using Lemma 3.7, assume T is large enough to ensure that

$$U(u)(t,y) \le \frac{a+b}{2}$$
 for $y_n \ge g(\hat{y}) + \frac{\delta}{2}$, $t \ge T$.

We can use the supersolution from the previous lemma for $t \geq T$ with g replaced by $h_{\tau(T)}$, δ replaced by $\delta/4$, $m = \varepsilon \theta(b-a)(\sqrt{|l|^2+1})/2 \equiv c_1 \varepsilon$, and

$$w(T,y) = \begin{cases} b & \text{for } y_n \leq h_{\tau(T)}(\hat{y}) - \delta/4, \\ (a+b)/2 & \text{for } h_{\tau(T)}(\hat{y}) - \delta/4 \leq y_n \leq h_{\tau(T)}(\hat{y}), \\ a & \text{for } h_{\tau(T)}(\hat{y}) < y_n. \end{cases}$$

It follows that for $t \geq T + \delta/4(\eta + c_1\varepsilon) \equiv T + c_2\delta$, $U(u)(t,y) \equiv a$ on the set $\{y: y_n > h_{\tau(T)}(\hat{y}) - c_1c_2\delta\varepsilon\}$. Thus $\tau_0 \leq \tau(T) - c_1c_2\delta\varepsilon$ which implies that $c_1c_2\delta\varepsilon \leq \varepsilon^2$. This is a contradiction since ε can be taken arbitrarily small. The proof of (1) is complete; (2) follows in a similar way using subsolutions.

An immediate consequence of this theorem is

THEOREM 3.12. Suppose u_0 is admissible, piecewise continuous, and k = [f(b) - f(a)]/(b-a). Then for any $\varepsilon > 0$ there is a constant $T_{\varepsilon} < \infty$ such that

$$u(t,x) = v(x - kt)$$
 for a.e. $x \in \mathbb{R}^n$,

when $t > T_{\varepsilon}$ and $\operatorname{dist}(x - kt, \partial M) > \varepsilon$.

4. The lower dimensional case. We now analyze the situation when d < n. By an appropriate rotation of the spacial coordinates we can assume that

$$f(c) = (h_1(c), \dots, h_d(c), j_1(c), \dots, j_{n-d}(c)),$$

dim span $\{\tilde{h}(c): a \leq c \leq b\} = d$, and $\tilde{j}(c) = 0$ for $a \leq c \leq b$, where

$$ilde{h}(c) \equiv h(c) - h(a) - (c-a)[h(b) - h(a)]/(b-a), \ ilde{j}(c) \equiv j(c) - j(a) - (c-a)[j(b) - j(a)]/(b-a).$$

Moreover the projection of μ on \mathbf{R}^d is nonzero and if $\mu' = P_{\mathbf{R}^d}(\mu)/||P_{\mathbf{R}^d}(\mu)||$ we have

$$\langle \mu', \tilde{h}(c) \rangle \le \langle \mu, \tilde{f}(c) \rangle \le -\theta(c-a)(b-c)$$
 for $a \le c \le b$.

We use the notation x = (x', x'') where $x' \in \mathbf{R}^d$ and $x'' \in \mathbf{R}^{n-d}$.

LEMMA 4.1. Suppose u_0 is admissible and k=(k',k'')=[f(b)-f(a)]/(b-a). Then for $t\geq T$ (where T is such that $a\leq u(t,x)\leq b$ for $t\geq T$) we have $u(t,x)=\tilde{u}(t,x'-k't;x''-k''t)$, where $\tilde{u}(t,x';x'')\in C([T,\infty);L^1_{\mathrm{loc}}(\mathbf{R}^d))$ for L^{n-d} a.e. x'' and is the solution of the Cauchy problem

(4.2)
$$\begin{cases} \partial_t w + \operatorname{div}_{x'} \tilde{h}(w) = 0 & \text{for } t \ge T, \ x' \in \mathbf{R}^d, \\ w(T, x') = u(T, x' + k'T; x'' + k''T). \end{cases}$$

Moreover for each such x'', u(T, x' + k'T; x'' + k''T) is admissible relative to a nondegenerate shock of the form

$$\varphi(x';x'') = \begin{cases} a & for \ \langle x' - x'_0(x''), \mu' \rangle > 0, \\ b & for \ \langle x' - x'_0(x''), \mu' \rangle < 0. \end{cases}$$

PROOF. As in §2 we assume that f(a)=f(b)=0 so that $\tilde{h}(c)=h(c),\,j(c)=0$ for $a\leq c\leq b,$ and k=0.

It suffices to prove the lemma in the region $\Omega_N = \{(t, x): T \leq t < T + N, |x| \leq N\}$ for any $N < \infty$. From (1.3) we see that u is uniquely determined on Ω_N by u(T, x)

with $|x| \le N(1+l)$. Thus without loss of generality we can assume that the solution is redefined at time t = T for |x| > N(1+l) so that $u(T, x) \equiv a$ for $|x| \ge 2N(1+l)$.

To begin with suppose $u(T,x) \in C^{\infty}(\mathbb{R}^n)$ and $h \in C^{\infty}(\mathbb{R}^1;\mathbb{R}^d)$. For $\varepsilon > 0$ let $v_{\varepsilon}(t,x';x'')$ be the solution to

$$\begin{cases} \partial_t v_{\varepsilon} + \operatorname{div}_{x'} h(v_{\varepsilon}) = \varepsilon \Delta_{x'} v_{\varepsilon} & \text{for } t \geq T \text{ and } x \in \mathbf{R}^n, \\ v_{\varepsilon}(T, x'; x'') = u(T, x) & \text{for } x \in \mathbf{R}^n. \end{cases}$$

Thus we view x'' as a parameter in the initial conditions. The theory of parabolic equations implies that a unique solution exists with $a \le v_{\varepsilon} \le b$ and

$$v_{\varepsilon} \in C^{\infty}([T,\infty) \times \mathbf{R}^n).$$

From [8, §§17.2 and 18.1], we obtain the estimates

(a)
$$\int_{\mathbf{R}^d} |\nabla_{x'} v_{\varepsilon}|(t, x'; x'') dx'$$

$$\leq \int_{\mathbf{R}^d} |\nabla_x v_{\varepsilon}|(T, x'; x'') dx' \quad \text{for } t \geq T, \ x'' \in \mathbf{R}^{n-d},$$

(b)
$$\int_{T}^{T'} \int_{\mathbf{R}^{d}} |\partial_{t} v_{\varepsilon}|(t, x'; x'') dx' dt \\ \leq c(T', l) \cdot \int_{\mathbf{R}^{d}} (|\nabla_{x'} v_{\varepsilon}|(T, x'; x'') + \varepsilon |\Delta_{x'} v_{\varepsilon}|(T, x'; x'')) dx'$$

for $T \leq T' < \infty$ and $x'' \in \mathbb{R}^{n-d}$. As a result we get

(c)
$$\int_{\mathbf{R}^n} |\nabla_x v_{\varepsilon}|(t,x) dx \leq \int_{\mathbf{R}^n} |\nabla_x v_{\varepsilon}|(T,x) dx \quad \text{for } t \geq T,$$

(d)
$$\int_{T}^{T'} \int_{\mathbf{R}^{n}} |\partial_{t} v_{\varepsilon}|(t, x) dx dt \\ \leq C(T', l) \cdot \int_{\mathbf{R}^{n}} (|\nabla_{x} v_{\varepsilon}|(T, x) + \varepsilon |\Delta_{x'} v_{\varepsilon}|(T, x)) dx.$$

If we assume only that $h \in C^1([a,b]; \mathbf{R}^d)$ then by taking a sequence of smooth functions, h_j , converging to h in this space one finds that the corresponding solutions, $v_{\varepsilon,j}$, together with $\nabla_x v_{\varepsilon,j}$ converge uniformly in $[T,T'] \times \mathbf{R}^n$. Also

$$\sup_{[T,T']\times R^n} |\partial_t v_{\varepsilon,j}| \le C < \infty$$

independent of j. Letting $j \to \infty$ the limit, v_{ε} , will be a solution of (4.3); the integrands in (a)-(d) are all locally integrable functions and since the right-hand sides of the four inequalities are independent of j they remain valid in the limit.

From (c) and (d) we see that a sequence $\varepsilon_j \downarrow 0$ can be chosen so that $v_{\varepsilon_j} \to w \in BV((T,\infty) \times \mathbb{R}^n)$ with convergence in $L^1_{loc}((T,\infty) \times \mathbb{R}^n)$. We can further assume that there is a set, Z, of full \mathcal{L}^{n-d} measure so that

$$v_{\varepsilon_i}(t,x) \to w(t,x';x'')$$
 in $L^1_{\text{loc}}((T,\infty) \times \mathbf{R}^d)$

for $x'' \in Z$. From (a) and (b) we see that $w(t, x'; x'') \in BV((T, \infty) \times \mathbf{R}^d)$ for each $x'' \in Z$. Thus (as in §1) we can define

$$\bar{w}(t,x';x'') \equiv \lim_{\tau \to \infty} \int_{B_{\tau}((t,x'))} w(\tau,z;x'') d\tau dz, \qquad t > T, \ x'' \in Z.$$

From [8, §18], for each such x'', $\bar{w}(t,x';x'')$ is the solution to (4.2). By the same argument as in [8, §18], $\bar{w}(t,x)$ (the Lebesgue limit in \mathbf{R}^{n+1}) is the solution to (1.1) with $\bar{w}(t,x) \to u(T,x)$ as $t \downarrow T$ in $L^1_{\mathrm{loc}}(\mathbf{R}^n)$. We point out that $\bar{w}(t,x';x'')$ agrees with $\bar{w}(t,x)$ as an element in $C([T,\infty);L^1_{\mathrm{loc}}(\mathbf{R}^n))$. We point out that $\bar{w}(t,x';x'')$ agrees with $\bar{w}(t,x)$ as an element in $C([T,\infty);L^1_{\mathrm{loc}}(\mathbf{R}^n))$. Thus the proof is complete when $u(T,x) \in C^{\infty}(\mathbf{R}^n)$.

For the general case consider a sequence of functions $u_j(T,x) \in C^\infty(\mathbf{R}^n)$ with $a \leq u_j \leq b, \ u_j(T,x) = a$ for $|x| \geq 2N(1+l)$, and $u_j \to u$ in $L^1_{\mathrm{loc}}(\mathbf{R}^n)$. Using equation (1.3) we see that $\{\bar{u}_j(t,x';x'')\}, \ j=1,2,\ldots$, forms a Cauchy sequence in $C([T,T+N];L^1_{\mathrm{loc}}(\mathbf{R}^n))$, and that for a.e. $x'' \in \mathbf{R}^{n-d}$, the sequence is Cauchy in $C([T,T+N];L^1_{\mathrm{loc}}(\mathbf{R}^d))$. For a subsequence, $\bar{u}_j(t,x';x'') \to u(t,x)$ for all $t \in [T,T+N]$ and a.e. $x'' \in \mathbf{R}^{n-d}$, pointwise almost everywhere in \mathbf{R}^d . This limit has the required properties.

With this result and those from §§2 and 3 we have the following.

THEOREM 4.4. Let $u_0(x)$ be admissible and let $b \cdot \chi_M(x-kt) + a \cdot \chi_{\mathbf{R}^n-M}(x-kt)$ be the asymptotic limit of u(t,x). Then for almost every $x_0'' \in \mathbf{R}^{n-d}$, $M(x_0'') \equiv \{x': (x',x_0'') \in M\}$ is a Lipschitz domain in \mathbf{R}^d . Moreover there is a rotation of the x' coordinates to $y' = (\hat{y},y_d) \in \mathbf{R}^{d-1} \times \mathbf{R}^1$ and a function $g(\hat{y},x'')$ defined for almost every x'', Lipschitz continuous in \hat{y} (with Lipschitz constant independent of x''), so that

$$M(x'') = \{ y' \in \mathbf{R}^d : y_d < g(\hat{y}, x'') \}.$$

LEMMA 4.5. Suppose $u_0(x)$ is admissible and piecewise continuous. Let \tilde{u} be as in Lemma 4.1 but with x' replaced by the rotated coordinates, $y'=(\hat{y},y_d)$. Then for almost every $x''\in \mathbf{R}^{n-d}$ and any $\delta>0$ there is a constant $T(\delta,x'')<\infty$ so that

$$\tilde{u}(t, y'; x'') = a$$
 for a.e. y such that $y_n \ge g(\hat{y}, x'') + \delta$, $t \ge T$, $\tilde{u}(t, y'; x'') = b$ for a.e. y such that $y_n < g(\hat{y}, x'') - \delta$, $t > T$.

PROOF. Assume that f(a) = f(b) = 0, so that k = 0 and $u = \tilde{u}$. Consider for each $j = 1, 2, \ldots, q_{1/j}$ and $p_{1/j}$ as defined in Lemma 3.7. For $t \geq t^*$, using Theorem 2.3,

$$a \leq q_{1/j}(t,x) \leq p_{1/j}(t,x) \leq b.$$

Now there exists a set Z of full \mathcal{L}^{n-d} measure so that for all j,

$$q_{1/j}(t^*, x', x'') \le u(t^*, x' + h, x'') \le p_{1/j}(t^*, x', x'')$$

for \mathcal{L}^d a.e. x', for each $|h| < \tau(1/j)$ and $x'' \in Z$. By construction $p_{1/j}(t^*, x) - q_{1/j}(t^*, x) \to 0$ as $j \to \infty$ in $L^1_{\text{loc}}(\mathbf{R}^n)$. Hence for almost every $x'' \in Z$,

$$\int_{|x'| \le R} [p_{1/j}(t^*, x', x'') - q_{1/j}(t^*, x', x'')] dx' \to 0$$

as $j \to \infty$ for any $R < \infty$. Using Lemma 4.1 we can argue as in Lemma 3.7 and Theorem 3.11 for each such x'' to obtain the conclusion of the lemma.

LEMMA 4.6. Let $u_0(x)$ be admissible and piecewise continuous. Then for all $\varepsilon > 0$ and almost every $z'' \in \mathbb{R}^{n-d}$ there is a constant, $T(\varepsilon, z'')$, such that if

$$\begin{split} x''-k''t&=z'',\\ u(t,x',x'')&=\tilde{u}(t,x'-k't;x''-k''t)\\ &=b\cdot\chi_{M(x''-k''t)}(x'-k't)+a\cdot\chi_{R^d-M(x''-k''t)}(x'-k't)\\ for\ t\geq T\ and\ almost\ every\ x'\ such\ that\ \mathrm{dist}(x'-k't,\partial M(x''-k''t))>\varepsilon. \end{split}$$

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